

Variational Iteration Sumudu Transform Method for Solving Fractional Nonlinear Gas Dynamics Equation

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Abstract: In this paper, generalized nonlinear fractional gas dynamics equation has been solved by using fractional variational iteration Sumudu transform method (VISTM). This technique was extended to derive analytical solutions in the form of a series for this equation which converges to the exact solution of the problem. It is indicated that the solutions obtained by the (VISTM) are reliable, very efficient, and professional techniques for solving different kinds of nonlinear fractional differential equation.

Keywords: Fractional calculus, Sumudu transform, Fractional variational iteration method, Variational iteration Sumudu transform method (VISTM), Nonlinear fractional gas dynamics equation.

1. INTRODUCTION

Nonlinear partial differential equations are useful in describing the various Phenomena in disciplines. The variational iteration method was proposed by He [1] and was successfully applied to autonomous ordinary differential equation [2] to nonlinear partial equations with variable coefficients [3] to Schrodinger-KdV, generalized KdV and shallow water equations [4] to linear Helmholtz partial differential equation [5] recently to nonlinear fractional differential equations with Caputo differential derivative [6, 7] and to other fields, [8]. The variational iteration method gives rapidly convergent successive approximations of the exact solution if such a solution exists; otherwise a few approximations can be used for numerical purposes. The method is effectively used in [2-4, 9-11] and the references there in Jafari et al [9, 10] applied the variational iteration method to the gas dynamics equation and Stefan problem. The objective of this paper is to extend the application of the fractional Variational iteration method combined with Sumudu transform method (VISTM) to solve nonlinear fractional gas dynamics equation.

2. FRACTIONAL CALCULUS

The fractional calculus has found diverse applications in various scientific and technological fields [12, 13], such as thermal engineering, acoustics, electromagnetism, control, robotics, viscoelasticity, diffusion, edge detection, turbulence, signal processing, and many other physical processes. Fractional differential equations (FDEs) have also been applied in modeling many physical, engineering problems, and fractional differential equations in nonlinear dynamics [14, 15]. In this section, we mention the following basic definitions of fractional calculus which are used further in the present paper:

Definition 1: The Riemann-Liouville fractional integral operator of order $\alpha > 0$, of a function $f(t) \in c_n$ and $p \ge -1$ is defined as [16].

$$I^{\alpha}f(t) = \frac{1}{\Gamma(\alpha)} \int_{0}^{t} (t-\tau)^{\alpha-1} f(\tau) d\tau, (\alpha > 0), I^{0}f(t) = f(t).$$

For the Riemann- Louville fractional integral we have:

$$I^{\alpha}t^{\gamma} = \frac{\Gamma(\gamma+1)}{\Gamma(\alpha+\gamma+1)}t^{\alpha+\gamma}$$

Definition 2: The Sumudu transform of the Caputo fractional derivative is defined as follows [17]:

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$$S[D_t^{\alpha} f(t)] = u^{-\alpha} S[f(t)] - \sum_{k=0}^{m-1} u^{-\alpha+k} f^{(k)}(0+), (m-1 < \alpha \le m)$$

Definition 3: The Mittage-Leffler function $E_{\alpha}(z)$ with $\alpha > 0$ is defined by the following series representation, valid in the whole complex plane [18]:

$$E_{\alpha}(z) = \sum_{k=0}^{\infty} \frac{z^{k}}{\Gamma(\alpha k + 1)}$$

Definition 4: Fractional derivative of compounded functions [19, 20] is defined as

$$d^{\alpha} f \cong \Gamma(\alpha + 1) df, 0 < \alpha < 1$$

Definition 5: The integral with respect to $(dx)^{\alpha}$ is defined as the solution of the fractional differential equation [19, 20]:

$$dy \cong f(x)(dx)^{\alpha}, x \ge 0, y(0) = 0, 0 < \alpha < 1$$

Definition 6: In early 90's, Watugala [21] introduced a new integral transform, named the Sumudu transform and applied it to the solution of ordinary differential equation in control engineering problems. The Sumudu transform is defined over the set of functions:

$$A = \left\{ f(t) : \exists M, \tau_1, \tau_2 > 0, | f(t) | < M e^{|t|/\tau_j}, \text{ if } t \in (-1)^j \times [0, \infty) \right\}$$

by the following formula

$$G(u) = S[f(t);u] = \int_{0}^{1} f(ut)e^{-t}dt, u \in (-\tau_{1}, \tau_{2}).$$

Some of the properties were established by Weerakoon in [22, 23]. In [24], by Asiru, further fundamental properties of this transform were also established. Similarly, this transform was applied to the one-dimensional neutron transport equation in [25] by Kadem. In fact it was shown that there is a strong relationship between Sumudu and other integral transforms; see Kılıcman et al. [26]. In particular the relation between Sumudu transform was extended to the distributions and some of their properties were also studied in Kılıcman and Eltayeb [28]. Recently, this transform is applied to solve the system of differential equations; see Kılıcman et al. in [29].

3. FRACTIONAL VARIATIONAL ITERATION METHOD (FVIM)

According to the Vim's rules [30 -32, 33] one need to follow three steps to determine the variational iteration formula:

- a. Establishing the correction functional,
- b. Identifying the Lagrange multipliers,
- c. Determining the initial iteration.

To describe the solution procedure of the fractional variational iteration method, we consider the following fractional differential equation:

$${}_{0}^{c}D_{t}^{\alpha}U + R[U(t)] + N[U(t)] = K(t), 0 < \alpha$$
⁽¹⁾

Where R is a linear operator, N is a nonlinear operator and K(t) is a given continuous function. Construct a correction functional:

$$U_{n+1} = U_n + \int_0^\tau \lambda(x,\tau) [{}_0^c D_\tau^\alpha U + R[U(\tau)] + N[U(\tau)] - K(\tau)] d\tau$$
(2)

Where $\lambda(x,\tau) = \frac{(-1)^{\alpha}}{\Gamma(\alpha)} (\tau - t)^{\alpha - 1}$ multiplier or a weighted function for any fractional order $\alpha > 0$,

$$U_{n+1} = U_n + \int_0^t \frac{(-1)^{\alpha}}{\Gamma(\alpha)} (\tau - t)^{\alpha - 1} [{}_0^c D_{\tau}^{\alpha} U + R[U(\tau)] + N[U(\tau)] - K(\tau)] d\tau$$
(3)

$$U_{n+1} = U_n - {}_0 I_t^{\alpha} [{}_0^{\alpha} D_t^{\alpha} U + R[U(t)] - K(t)]$$
(4)

The above formula is also valid for the FDEs with the R-L derivative.

It is obvious that the sequential approximations $U_k, k \ge 0$ can be established by determining λ , a general Lagrange's multiplier which can be identified optimally with the variational theory. The function U_n is a restricted variation which means $\partial U_n = 0$. Therefore, we first designate the Lagrange multiplier λ that will be identified optimally via integration by parts. The successive approximations $U_{n+1}(x,t), n \ge 0$ of the solution U(x,t), will be readily obtained upon using the obtained Lagrange multiplier and by using any selective function U_0 . The initial values are usually used for choosing the zeroth approximation U_0 . With λ determined, then several approximations, U_k , $k \ge 0$ follows immediately [34]. Consequently the exact solution may be procured by using:

$$U(x,t) = \lim_{n \to \infty} U_n(x,t)$$
⁽⁵⁾

4. FRACTIONAL VARIATIONAL ITERATION SUMUDU TRANSFORM METHOD

In the case of an algebraic equation f(x) = 0, the Lagrange multipliers can be evaluated by an iteration formula for finding the solution of the algebraic Equation f(x) = 0 that can be constructed as

$$x_{n+1} = x_n + \lambda f(x_n). \tag{6}$$

The optimality condition for the extreme $\frac{\partial x_{n+1}}{\partial x_n} = 0$ Leads to

$$\lambda = -\frac{1}{f'(x_n)},\tag{7}$$

Where ∂ is the classical variational operator. From (6) and (7), for a given initial value X_0 , we can find the approximate solution X_{n+1} by the iterative scheme for (6) as follows:

$$x_{n+1} = x_n - \frac{f(x_n)}{f'(x_n)}, f(x_0) \neq 0, n = 0, 1, 2, \dots$$
(8)

This algorithm is well known as the Newton-Raphson method and has quadratic convergence. Now, we extend this idea to finding the unknown Lagrange multiplier. The main step is to first take the Sumudu transform to Eq. (1). Then the linear part is transformed into an algebraic equation as follows:

$$\frac{G(u)}{u^{\alpha}} - \frac{U(0)}{u^{\alpha}} \dots \dots - \frac{U^{(n-1)}(0)}{u^{\alpha - (n-1)}} + S[R[U(t)] + N[U(t)] - K(t)] = 0,$$

$$G(u) = S[U(t)] = \int_{0}^{\infty} U(ut)e^{-t}dt, -\tau_{1} \prec u \prec \tau_{2}.$$
(9)
Where

The iteration formula of (8) can be used to suggest the main iterative scheme involving the Lagrange multiplier as:

$$G_{n+1} = G_n + \lambda \left(\frac{G(u)}{u^{\alpha}} - \frac{U(0)}{u^{\alpha}} \dots - \frac{U^{(n-1)}(0)}{u^{\alpha-(n-1)}} + S[R[U(t)] + N[U(t)] - K(t)]\right).$$
(10)

Considering S[R[U(t)] + N[U(t)] as restricted terms, one can derive a Lagrange multiplier as:

$$\partial G_{n+1} = \partial G_n + \partial (\lambda \frac{G_n(u)}{u^{\alpha}})$$

$$\partial G_{n+1} = \partial G_n + \frac{1}{u^{\alpha}} (\lambda' G_n + \lambda \partial G_n)$$

This yields the stationary conditions of Eq. (10) as follows;

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$$\partial G_n (1 + \frac{\lambda}{u^{\alpha}}) = 0,$$

$$\frac{1}{u^{\alpha}} (\lambda' G_n) = 0$$

$$\therefore \lambda = -u^{\alpha}$$

With Eq. (10) and the inverse-Sumudu transform S^{-1} , the iteration formula (9) can be explicitly given as:

$$U_{n+1} = U_n - S^{-1} \left[u^{\alpha} \left(\frac{G(u)}{u^{\alpha}} - \frac{U(0)}{u^{\alpha}} \dots - \frac{U^{(n-1)}(0)}{u^{\alpha-(n-1)}} + S[R[U(t)] + N[U(t)] - K(t)] \right) \right]$$
(11)

$$U_{n+1} = S^{-1} \left[u^{\alpha} \left(\frac{U(0)}{u^{\alpha}} \dots \dots \right) - \frac{U^{(n-1)}(0)}{u^{\alpha-(n-1)}} \right) - u^{\alpha} \left(S[R[U(t)] + N[U(t)] - K(t)) \right],$$
(12)

Consequently the exact solution may be procured by using

$$U(x,t) = \underset{n \to \infty}{\text{Lim}} U_n(x,t)$$
(13)

5. AN APPLICATION

In this paper, we consider the following nonlinear time- fractional gas dynamics equation of the form:

$$D_t^{\alpha}U + \frac{1}{2}(U^2)_x - U(1-U) = 0, t > 0, 0 < \alpha \le 1,$$
(14)

with the initial condition:

$$U(x,0) = e^{-x}$$
(15)

where α is parameter describing the order of the fractional derivative. The function U(x,t) is the probability density function, t is the time and x is the spatial coordinate. The derivative is understood in the Caputo sense. The general response expression contains parameter describing the order of the fractional derivative that can be varied to obtain various responses. In the case of $\alpha = 1$ the fractional gas dynamics equation reduces to the classical gas dynamics equation. The gas dynamics equations are based on the physical laws of conservation, namely, the laws of conservation of mass, conservation of momentum, conservation of energy etc. The nonlinear fractional gas dynamics has been studied previously by Das and Kumar [35].

Further, we apply the (VISTM) to solve the nonlinear time- fractional gas dynamics equation. The objective of the present paper is to extend the application of the (VISTM) to obtain analytic and approximate solutions to the time-fractional gas dynamics equation.

The Solution

After taking the Sumudu transform on both sides of Eq. (14) the iteration formula of Eq. (14) can be constructed as:

$$G_{n+1} = G_n + \lambda(u) \left[\frac{G_n}{u^{\alpha}} - \frac{\tilde{U}_0}{u^{\alpha}} + S\left[\frac{1}{2} (\tilde{U}_n^2)_x - \tilde{U}_n (1 - \tilde{U}_n) \right] \right]$$
(16)

where λ is a general Lagrange multiplier, which can be identified optimally via the variational theory, $\tilde{U}_0 = 0$ and $S[\frac{1}{2}(\tilde{U}_n^2)_x - \tilde{U}_n(1 - \tilde{U}_n)]$ is a restricted variation, that is, $\partial \tilde{U}_n = 0$.

i.e.

$$\partial G_{n+1} = \partial G_n + \partial G_n \frac{\lambda}{u^{\alpha}}$$
(17)

This yields the stationary conditions, which gives $\lambda = -u^{\alpha}$. Substituting this value of Lagrangian multiplier in (16) we get the following iteration formula:

$$G_{n+1} = G_n - u^{\alpha} \left[\frac{G_n}{u^{\alpha}} - \frac{U_0}{u^{\alpha}} + S\left[\frac{1}{2} (U_n^2)_x - U_n (1 - U_n) \right] \right].$$
(18)

Applying inverse Sumudu transform on both sides of Eq. (18) we get:

$$U_{n+1} = U_n - S^{-1} (u^{\alpha} [\frac{G_n}{u^{\alpha}} - \frac{U_0}{u^{\alpha}} + S[\frac{1}{2}(U_n^2)_x - U_n(1 - U_n)]])$$
(19)

with the initial condition

$$U(x,0) = e^{-x}$$

i.e.
$$U_0 = e^{-x} ,$$

$$U_1 = e^{-x} + e^{-x} \frac{t^{\alpha}}{\Gamma(\alpha+1)},$$

$$U_2 = e^{-x} + e^{-x} \frac{t^{\alpha}}{\Gamma(\alpha+1)} + e^{-x} \frac{t^{2\alpha}}{\Gamma(2\alpha+1)},$$

$$U_3 = e^{-x} + e^{-x} \frac{t^{\alpha}}{\Gamma(\alpha+1)} + e^{-x} \frac{t^{2\alpha}}{\Gamma(2\alpha+1)} + e^{-x} \frac{t^{3\alpha}}{\Gamma(3\alpha+1)}, \dots$$

Finally, we approximate the analytical solution $U(x,t)$ by:

$$U(x,t) = e^{-x} \sum_{k=0}^{\infty} \frac{(t^{\alpha})}{\Gamma(\alpha k+1)}$$

$$(20)$$

Where
$$\sum_{k=0}^{\infty} \frac{(t^{\alpha})^{\alpha}}{\Gamma(\alpha k+1)} = E_{\alpha}(t^{\alpha})$$
 is the famous Mittage–Leffler function i.e. Eq. (20) takes the form:

$$U(x,t) = e^{-x} E_{\alpha}(t^{\alpha})$$
(21)

In special case at $\alpha = 1$ we have

$$U(x,t) = e^{-x} \sum_{k=0}^{\infty} \frac{t^k}{\Gamma(k+1)} = e^{-x} e^t = e^{t-x}$$

$$\therefore U(x,t) = e^{t-x}.$$
(22)

Which is the same as given by HPM and HPSTM (see Jafari et al [36], and Singh et al [37]).

Now, we calculate numerical results of the probability density function U(x,t) for different timefractional Brownian motions $\alpha = 1/3, 1/2, 1$ and for various values of t and x. The numerical results for the approximate solution (21) obtained by using VIST M and the exact solution (22) for various values of t, x, and α are shown in Figure (a)-(d). It is observed from Figures (a) and (b) that U(x,t) increase with the increase in t and U(x,t) decrease with the increase in α . Figure 1(c) and figure 2 (d) clearly show that, when $\alpha = 1$, the approximate solution (21) obtained by the present method is very near to the exact solution. It is evident that the efficiency of the present method can be dramatically.



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Figure 1. The behavior of the U(x, t), x and t are obtained when (a) $\alpha = 0.75$, (b) $\alpha = 0.90$, (c) $\alpha = 1$.



Figure 1. The behavior of the U(x, t) and t are obtained at (---) α =0.90, (---), α =0.75(----), α =1.

6. CONCLUSION

In this paper, variational iteration Sumudu transform method (VISTM) is successfully applied for solving nonlinear time- fractional gas dynamics equation. This method is very powerful and efficient techniques for solving different kinds of linear and nonlinear fractional differential equations arising

in different fields of science and engineering. With the approach given in this paper, we can easily derive Lagrange multipliers without tedious calculation and new variational iteration formulae can be derived. Some FDEs with the Caputo derivatives are illustrated. The results show the modified method's efficiency compared with other versions of the VIM in fractional calculus. In conclusion, the VISTM may be considered as a nice refinement in existing numerical techniques and might find the wide applications.

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